Evolution of electrical current systems in the solar atmosphere: what are the key questions, and why?

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LWS workshop Boulder Colorado September 10-13 2007

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Background

Energy release in the corona

- Solar flares and CMEs involve release of magnetic energy
 - 'release of free energy associated with currents'



http://trace.lmsal.com/POD/



$$E = \frac{\mu_0}{8\pi} \int d^3x \int d^3x' \frac{\mathbf{J}(\mathbf{x}) \cdot \mathbf{J}(\mathbf{x}')}{|\mathbf{x} - \mathbf{x}'|}$$
(1)

 \blacktriangleright Current systems: $\boldsymbol{J}_{c},~\boldsymbol{J}_{p},~\boldsymbol{J}_{s}:$



- we don't know what changes during flares
- consequently the 'free energy' is not well defined

- Free energy is often defined as $E E_0$ where E_0 refers to the potential field **B**₀ matching the observed B_n (e.g. Low 1982)
 - ► Thomson's principle: B₀ is the minimum energy field for a fixed B_n
 - belief that photospheric B_n does not change during a flare
- Extremal energy state is a linear force-free field for fixed B_n and helicity (Woltjer 1958; 1959)
 - extremal energy state is a nonlinear force-free field when B_n is fixed and connectivity is specified (Sakurai 1989)
- However, flares may be driven by emerging flux, in which case B_n is changing? (e.g. Heyvaerts, Priest & Rust 1977; Zirin 1983)
- Can still estimate E, ΔE

Poynting theorem (Maxwell's equations):

$$\frac{\partial}{\partial t} \left(\frac{B^2}{2\mu_0} + \frac{\epsilon_0 E^2}{2} \right) + \nabla \cdot \left(\frac{\mathbf{E} \times \mathbf{B}}{\mu_0} \right) = -\mathbf{J} \cdot \mathbf{E}$$
(2)

J · E is the rate work is done per unit volume by the field
Magnetohydrodynamics (simple Ohm's law):

$$\frac{\partial}{\partial t} \left(\frac{1}{2} \rho u^2 + \frac{p}{\gamma - 1} \right) + \nabla \cdot \left(\frac{1}{2} \rho u^2 \mathbf{u} + \frac{\gamma}{\gamma - 1} \rho \mathbf{u} \right) - \mathbf{J} \cdot \mathbf{E}$$
$$= -\frac{J^2}{\sigma}$$
(3)

•
$$-J^2/\sigma$$
 describes dissipation

• in MHD σ hides all sins

Observations

Photospheric fields

- Spectro-polarimetric measurements of photospheric lines permit inference of the vector magnetic field
 - it's an inference... caveat emptor
- New instruments (Hinode/SOT, SDO/HMI, SOLIS)



- Inversion problem: Stokes profiles \rightarrow **B**_p
 - least squares fitting to an analytic radiative transfer solution (e.g. Unno 1956; Rachovsky 1962; Auer, Heasley & House 1977; Skumanich & Lites 1987; Jefferies, Lites & Skumanich 1989)
 - model assumes Milne-Eddington atmosphere, constant physical quantities
 - observed profiles often do not have symmetries of the model
 - weak-field approximations used also (e.g. Jefferies & Mickey 1991)
 - More recent method: Stokes Inversion based on Response functions (SIR) (Ruiz Cobo & del Toro Iniesta 1992; Westendorp Plaza et al. 1998; 2001)
- \blacktriangleright 180 degree ambiguity in direction of ${\bm B}_{\rm p\perp}$
 - various methods of resolution (e.g. Metcalf et al. 2006)
- Error in spectro-polarimetric measurements

Photospheric currents

Vertical component of current density (Moreton & Severny 1968)

$$J_{\mathrm{p}z} = \frac{1}{\mu_0} \left(\frac{\partial B_{\mathrm{p}y}}{\partial x} - \frac{\partial B_{\mathrm{p}x}}{\partial y} \right) \tag{4}$$

- finite differencing of vector magnetic field values
- substantial random errors (e.g. Leka & Skumanich 1999)
- harder to categorise systematic uncertainty
- some authors very skeptical (e.g. Parker 1996; cf. McClymont, Jiao & Mikic 1997; Semel & Skumanich 1998)
- however, essentially the only quantitative information on current systems (vertical component of J_p)
- Components J_{px} , J_{py} not available ($\partial \mathbf{B}/\partial z$ required)

Properties of currents

- ▶ Maximum values $J_{\mathrm{p}z} \approx 10-50\,\mathrm{mA}\,\mathrm{m}^{-2}$ (Gary & Démoulin 1995; Leka & Skumanich 1999)
- Persist for ≈ 1 day (Pevtsov, Canfield, Metcalf 1994; Schrijver et al. 2005)
- Photospheric J_{pz} qualitatively consistent with coronal twist (Pevtsov, Canfield & McClymont 1997)
- Origin of photospheric currents is unknown
 - subphotospheric: flux emerges with currents flowing ('pre-stressed') (e.g. McClymont & Fisher 1989; Leka et al. 1996)
 - photospheric: coronal fields twisted by photospheric flows
- Photospheric model implies 'neutralized' patterns of current flow in a given polarity (e.g. Melrose 1991; Aulanier, Démoulin & Grappin 2005)
 - observations ambiguous (e.g. Wilkinson et al. 1992; Wheatland 2000)
 - Hinode data suggests currents are neutralized





- Current helicity $H_{\rm c} = \int d^3x \, {f B} \cdot {f J}$
 - ► in common with magnetic helicity $(H_m = \int d^3 x \mathbf{A} \cdot \mathbf{B})$ a measure of twist
 - ▶ component of the integrand (B_{pz}J_{pz}) determined at photosphere
 - related to force-free parameter $\alpha_{\rm p}=\mu_0 J_{{\rm p}z}/B_{{\rm p}z}$
- Hemispheric preference for sign of current helicity (e.g. Richardson 1941; Seehafer 1990; Rust 1994)
 - left handed features in north, right-handed in south
 - however, localised variation in e.g. sign of $\alpha_{
 m p}$
- Mechanism responsible remains obscure (e.g. van Driel-Gesztelyi, Démoulin, & Mandrini 2003)

Modelling

Force-free fields

Zeroth order static model

$$(\nabla \times \mathbf{B}) \times \mathbf{B} = 0 \qquad \nabla \cdot \mathbf{B} = 0$$
 (5)

or

$$\nabla \times \mathbf{B} = \alpha \mathbf{B} \qquad \mathbf{B} \cdot \nabla \alpha = \mathbf{0} \tag{6}$$

• spatially varying α (nonlinear model) necessary

• Dependent variable **B** or else **B**, α

- BCs: B_n and α on one sign of B_n
- $\bullet \ \alpha_{\rm p} = \mu_0 J_{\rm pz} / B_{\rm pz}$

Elliptic/hyperbolic PDEs may be numerically solved

- current-field iteration (Grad & Rubin 1958)
- magnetofrictional method (Chodura & Schlueter 1981)
- vertical integration (Wu et al. 1990)
- optimization method (Wheatland, Sturrock, Roumeliotis 2000)
- fast current-field iteration (Wheatland 2006; 2007)

- \blacktriangleright Application to \boldsymbol{B}_{p} difficult
 - photospheric field is not force free (e.g. Metcalf et al. 1995) preprocessing? (Wiegelmann, Inhester & Sakurai 2006)
 - difficulty of determining \(\alpha_p\)
- Goals: estimate coronal energy, topology
 - MHD virial theorem provides energy from B_p:

$$E = \frac{1}{\mu_0} \int_{\mathbf{p}} (B_{\mathbf{p}x} x + B_{\mathbf{p}y} y) B_{\mathbf{p}z} dx dy \tag{7}$$

(Chandrasekhar 1961; Moldensky 1974)

- assumes force-free BCs
- NLFFF workshops
 - 2005: Low & Lou test cases (Schrijver et al. 2006)
 - 2006: van Ballegooijien field (Metcalf et al. 2007)
 - 2007: Hinode/SOT AR 10930 (Metcalf/DeRosa, LWS 2007)



Metcalf et al. 2007

- Other limitations/problems:
 - difficulty of solution of PDEs
 - topologically significant sites not force-free?
 - sequence of force-free states is not equivalent to MHD evolution (e.g. Syrovatskii 1978)

Magnetohydrostatics (MHS)

MHS is next simplest model

$$abla p = \mathbf{J} \times \mathbf{B} \qquad
abla \times \mathbf{B} = \mu_0 \mathbf{J} \qquad
abla \cdot \mathbf{B} = 0 \qquad (8)$$

p is constant on lines of **B** and **J**

▶ Writing $\mathbf{J} = \mathbf{J}_{\parallel} + \mathbf{J}_{\perp}$ with $\mathbf{J}_{\parallel} = \alpha \mathbf{B} / \mu_0$

$$\mathbf{J}_{\perp} = \frac{1}{B^2} \mathbf{B} \times \nabla p \quad \text{and} \quad \mathbf{B} \cdot \nabla \alpha + \mu_0 \nabla \cdot \mathbf{J}_{\perp} = 0$$
 (9)

▶ BCs: B_n plus p, J_n on one sign of B_n (Grad & Rubin 1958)

May be solved by Grad-Rubin iteration

- ► start with **B**₀
- propagate boundary values of p along field lines
- determine J_{\perp} everywhere
- Solve for α along field lines given J_n and \mathbf{J}_{\perp}
- Obtain a new **B** by solving $\nabla \times \mathbf{B} = \mu_0 \mathbf{J}$ with $\nabla \cdot \mathbf{B} = 0$
- iterate...

- Can be generalised to include gravity
 - ρ enters the force balance
 - p varies along field lines
 - equations (8) supplemented by EOS, assumption about T
- In principle could be applied to photospheric data
 - B_n , J_n from \mathbf{B}_p
 - photospheric pressure diagnostic? (p_e from SIR?)
 - accommodates non force-free BCs
- Grad-Rubin iteration is a generalisation of current-field iteration (e.g. Wheatland 2006)
 - not significantly more computationally intensive

Magnetohydrodynamics (MHD)

- Minimum model for 'evolution of electrical current systems'
 - MHD modelling to date highly idealised
 - rarely applied to data (cf. Riley, Linker, Mikic 2001; Peter, Gudikson & Nordlund 2004; W. Abbett LWS 2007)
- Dependent variables **v**, **B**, p, ρ
 - $\mathbf{J} = \mu_0^{-1} \nabla \times \mathbf{B}$ considered a secondary quantity (e.g. Parker 1996)
 - simplistic BCs on v, B may imply unrealistic currents at boundaries, e.g. 'line tying' (Melrose 1991; 1995)
- Time dependent BCs
 - ▶ $\mathbf{B}_{\mathrm{p}}(t)$ available; $\mathbf{v}_{\mathrm{p}}(t)$, $\rho_{\mathrm{p}}(t)$, $\rho_{\mathrm{p}}(t)$ less accessible
 - correlation tracking for $\mathbf{v}_{\mathrm{p}}(t)$? (e.g. November & Simon 1988)
 - photospheric density/pressure diagnostics? (p_e(t) from SIR?)
 - is it possible to perform real-time MHD modelling from data?
- Numerical solution of equations remains challenging
 - hyperbolic/parabolic PDEs

Conclusions

Summary

- Coronal currents determine magnetic energy and dissipation
- Observing current systems is problematic
 - J_{pz} is inferred with substantial uncertainty
 - nature of currents not well understood
 - new instruments should provide substantially improved data
- Limited ability to model coronal fields, currents from BCs
 - force-free, MHS and MHD models described
 - prospects for application to data discussed

Questions

- Can we better define 'free energy'?
- Can we improve Stokes inversion?
 - can we extract additional information, e.g. p_e?
- Can we confirm the reality of inferred values of J_{pz}?
- Can we identify the origin of the currents?
- Can we perform coronal field modelling from BCs?
 - is force-free modelling useful/possible?
 - would magnetostatic modelling be better?
 - is time-dependent MHD modelling possible?